

# マルチセル円筒殻の応力解析について

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## 1. まえがき

著者らは、先に、マルチセル円筒殻に鉛直荷重を中央に載荷した場合の応力解析を行なった<sup>(1)</sup>。本報告では、荷重状態として、図1に示す場合を考え、この荷重状態の時のマルチセル円筒殻の応力解析を行う。解析では、長軸方向に、Finite Sine Transformationsを、断面回転方向に、Finite Fourier Transformationsを行う。  
このFinite Fourier Transformationsを行う事によって、単に、8元の連立方程式を解く問題に帰着する。

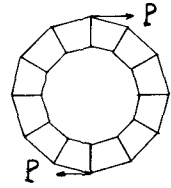


図.1

## 2. 析板要素の断面力

図2に示す細長い矩形板要素において、平面応力状態から導き出される剪断力 $T$ 、法線力 $S$ 、及び、平板の曲げの問題から導き出される曲げモーメント $M$ 、剪断力 $X$ は、各々、次の如くなる。

$$T_{j,j+1} = \frac{N\dot{u}_j}{\delta}(2\dot{u}_j + \dot{u}_{j+1}) + \frac{N\dot{u}_j}{2}(\dot{u}_{j+1} - \dot{u}_{j,j+1}) + \frac{Gt}{2}(\dot{u}_{j,j+1} + \dot{u}_{j,j+1}) + \frac{Gt}{\delta}(u_{j+1} - u_j)$$

$$T_{j,j-1} = \frac{N\dot{u}_j}{\delta}(2\dot{u}_j + \dot{u}_{j-1}) + \frac{N\dot{u}_j}{2}(\dot{u}_{j-1} - \dot{u}_{j,j-1}) - \frac{Gt}{2}(\dot{u}_{j,j-1} + \dot{u}_{j,j-1}) - \frac{Gt}{\delta}(u_j - u_{j-1})$$

$$S_{j,j+1} = \frac{N}{\delta}(u_{j,j+1} - u_{j,j+1}) + \frac{N\dot{u}_j}{2}(\dot{u}_j + \dot{u}_{j+1}) + \frac{Gt}{2}(\dot{u}_{j,j+1} - \dot{u}_j) + \frac{GtR}{\delta}(2\dot{u}_{j,j+1} + \dot{u}_{j,j+1})$$

$$S_{j,j-1} = \frac{N}{\delta}(u_{j,j-1} - u_{j,j-1}) + \frac{N\dot{u}_j}{2}(\dot{u}_{j-1} + \dot{u}_j) - \frac{Gt}{2}(\dot{u}_j - \dot{u}_{j-1}) - \frac{GtR}{\delta}(2\dot{u}_{j,j-1} + \dot{u}_{j,j-1})$$

$$M_{j,j+1} = \frac{2K_l}{\delta}\{2\phi_j + \phi_{j+1} - \frac{\phi}{\delta}(u_{j,j+1} - u_{j,j+1})\} - \nu K_l \ddot{u}_{j,j+1}$$

$$M_{j,j-1} = \frac{2K_l}{\delta}\{2\phi_j + \phi_{j-1} - \frac{\phi}{\delta}(u_{j,j-1} - u_{j,j-1})\} - \nu K_l \ddot{u}_{j,j-1}$$

$$X_{j,j+1} = -\frac{\delta K_l}{\delta^2}(\phi_j + \phi_{j+1}) + K_l \ddot{\phi}_j - \frac{12K_l}{\delta^3}(u_{j,j+1} - u_{j,j+1}) + \frac{2K_l}{\delta}(\ddot{u}_{j,j+1} - \ddot{u}_{j,j+1}) - \frac{K_l R}{\delta}(2\ddot{u}_{j,j+1} + \ddot{u}_{j,j+1})$$

$$X_{j,j-1} = -\frac{\delta K_l}{\delta^2}(\phi_j + \phi_{j-1}) + K_l \ddot{\phi}_j - \frac{12K_l}{\delta^3}(u_{j,j-1} - u_{j,j-1}) + \frac{2K_l}{\delta}(\ddot{u}_{j,j-1} - \ddot{u}_{j,j-1}) + \frac{K_l R}{\delta}(2\ddot{u}_{j,j-1} + \ddot{u}_{j,j-1})$$

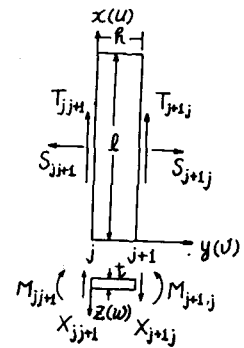


図.2

ここで、 $\nu$ ; ポアソン比,  $N = Et/(1-\nu^2)$ ,  $K = Et^3/12(1-\nu^2)$ ,  $(\cdot)$ ;  $x$ に関する微分

## 3. 解析

i) 節点における力のつりあい

図3に示す如く、外側と内側の各節点での力のつりあいは、

$$1) M_{j,j-1} - M_{j,j+1} - M_{j,j} = 0$$

$$2) T_{j,j-1} + T_{j,j+1} + T_{j,j} = 0$$

$$3) P_j + S_{j,j+1} \cos \alpha - S_{j,j-1} \cos \alpha - X_{j,j+1} \sin \alpha - X_{j,j-1} \sin \alpha - X_{j,j} = 0$$

$$4) S_{j,j+1} \sin \alpha + S_{j,j-1} \sin \alpha + X_{j,j+1} \cos \alpha - X_{j,j-1} \cos \alpha + S_{j,j} = 0$$

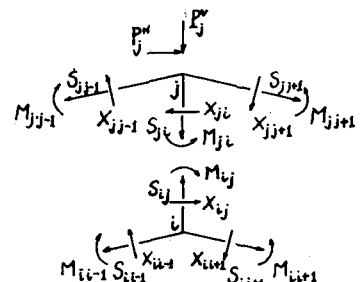


図.3

$$5) M_{i,i-1} + M_{ij} - M_{i,i+1} = 0$$

$$6) T_{i,i-1} + T_{ij} - T_{i,i+1} = 0$$

$$7) S_{i,i-1} \cos \alpha - S_{i,i+1} \cos \alpha - X_{i,i+1} \sin \alpha - X_{i,i-1} \sin \alpha + X_{ij} = 0$$

$$8) S_{i,i+1} \sin \alpha + S_{i,i-1} \sin \alpha + X_{i,i+1} \cos \alpha - X_{i,i-1} \cos \alpha - S_{ij} = 0$$

以上、1) ~ 8) 式に、2) 式示した各断面力を代入する。次に、これを Infinite Fourier Transformation を行い、 $\alpha = 0, \pi$  において、その断面では無限剛性で有し、面外には自由に動けるサイヤラムを有し、かつ単純支持であると仮定する。

又、変形は、次に示す如く半経方向と接線方向の変形に変換したものをを用いる。

$$U_{j,i+1} = U_j^0 \cos \alpha + W_j^0 \sin \alpha \quad U_{j,i-1} = U_j^0 \cos \alpha - W_j^0 \sin \alpha$$

$$U_{j,i+1} = U_{j,i}^0 \cos \alpha - W_{j,i}^0 \sin \alpha \quad U_{j,i-1} = U_{j,i}^0 \cos \alpha + W_{j,i}^0 \sin \alpha$$

$$W_{j,i+1} = -U_j^0 \sin \alpha + W_j^0 \cos \alpha \quad W_{j,i-1} = U_j^0 \sin \alpha + W_j^0 \cos \alpha$$

$$W_{j,i+1} = U_{j,i}^0 \sin \alpha + W_{j,i}^0 \cos \alpha \quad W_{j,i-1} = -U_{j,i}^0 \sin \alpha + W_{j,i}^0 \cos \alpha$$

$$U_{j,i} = W_j^0 \quad W_{j,i} = -U_j^0$$

内側断面についてでも同様である。これによって 1) ~ 8) 式を整理解ると

$$1)' \quad 4(2a_1 + a_2) \tilde{U}_j^0 + 2a_1 (\tilde{U}_{j+1}^0 + \tilde{U}_{j-1}^0) + 2a_2 \tilde{U}_i^0 - (2a_1 \sin \alpha + a_2) \tilde{U}_j^0 \\ - 6b_1 \sin \alpha (\tilde{U}_{j+1}^0 + \tilde{U}_{j-1}^0) - 6b_1 \cos \alpha (\tilde{W}_{j+1}^0 - \tilde{W}_{j-1}^0) = 0$$

$$2)' \quad (2\beta_1 + \beta_2) \tilde{U}_j^1 - \gamma_2 \tilde{U}_i^1 - \gamma_1 (\tilde{U}_{j+1}^1 + \tilde{U}_{j-1}^1) + \mu_2 \tilde{W}_j^1 - \varepsilon_2 \tilde{W}_j^0 - 2\varepsilon_1 \sin \alpha \tilde{W}_j^0 \\ - \mu_1 \sin \alpha (\tilde{W}_{j+1}^0 + \tilde{W}_{j-1}^0) + \mu_1 \cos (\tilde{U}_{j+1}^0 - \tilde{U}_{j-1}^0) = 0$$

$$3)' \quad -(2\varepsilon_1 \sin \alpha + \varepsilon_2) \tilde{U}_j^0 - 6b_1 \sin \alpha (\tilde{U}_{j+1}^0 + \tilde{U}_{j-1}^0) - 6b_2 \tilde{U}_i^0 + \left(\frac{L}{m\pi}\right)^2 \mu_1 \cos \alpha (\tilde{U}_{j+1}^0 - \tilde{U}_{j-1}^0) \\ + (2\xi_1 \cos^2 \alpha + 2\eta_1 \sin^2 \alpha + \eta_2) \tilde{U}_j^0 - \chi_2 \tilde{U}_j^1 + (\chi_1 \sin^2 \alpha - \xi_1 \cos^2 \alpha) (\tilde{U}_{j+1}^0 + \tilde{U}_{j-1}^0) \\ + (\chi_1 + \xi_1) \sin \alpha \cos \alpha (\tilde{W}_{j+1}^0 - \tilde{W}_{j-1}^0) = \tilde{F}_j$$

$$4)' \quad 6b_1 \cos \alpha (\tilde{U}_{j+1}^0 - \tilde{U}_{j-1}^0) - \left(\frac{L}{m\pi}\right)^2 (2\varepsilon_1 \sin \alpha + \varepsilon_2) \tilde{U}_j^1 - \mu_2 \left(\frac{L}{m\pi}\right)^2 \tilde{U}_i^1 - \xi_2 \tilde{W}_j^1 \\ - \mu_1 \left(\frac{L}{m\pi}\right)^2 \sin \alpha (\tilde{U}_{j+1}^1 + \tilde{U}_{j-1}^1) + (2\delta \sin^2 \alpha \xi_1 + 2\cos^2 \alpha \eta_1 + \xi_2) \tilde{W}_j^0 \\ - (\xi_1 + \chi_1) \sin \alpha \cos \alpha (\tilde{U}_{j+1}^0 - \tilde{U}_{j-1}^0) + (\sin^2 \alpha \xi_1 - \cos^2 \alpha \chi_1) (\tilde{U}_{j+1}^0 + \tilde{U}_{j-1}^0) = 0$$

$$5)' \quad -4(2a_3 + a_2) \tilde{U}_i^1 - 2a_3 (\tilde{U}_{i+1}^1 + \tilde{U}_{i-1}^1) - 2a_2 \tilde{U}_j^1 + (2a_3 \sin \alpha + a_2) \tilde{U}_j^1 \\ + 6b_3 \sin \alpha (\tilde{U}_{j+1}^1 + \tilde{U}_{j-1}^1) + 6b_3 \cos \alpha (\tilde{W}_{j+1}^1 - \tilde{W}_{j-1}^1) - 6b_2 \tilde{U}_j^0 = 0$$

$$6)' \quad -(2\beta_3 + \beta_2) \tilde{U}_i^1 + \gamma_3 (\tilde{U}_{i+1}^1 + \tilde{U}_{i-1}^1) + \gamma_2 \tilde{U}_j^1 - \mu_3 \cos \alpha (\tilde{U}_{j+1}^1 - \tilde{U}_{j-1}^1) \\ + \mu_3 \sin \alpha (\tilde{W}_{j+1}^1 + \tilde{W}_{j-1}^1) + (2\varepsilon_3 \sin \alpha - \varepsilon_2) \tilde{W}_j^1 + \mu_2 \tilde{W}_j^0 = 0$$

$$\begin{aligned} \eta) & \mu_3 \cos \alpha \left( \frac{1}{m\pi} \right)^2 (\tilde{U}_{i+1} - \tilde{U}_{i-1}) + \xi_3 \cos^2 \alpha (\tilde{V}_{j+1}^T + \tilde{V}_{j-1}^T) - \xi_3 \sin \alpha \cos \alpha (\tilde{\omega}_{j+1}^T - \tilde{\omega}_{j-1}^T) \\ & - (2\xi_3 \cos^2 \alpha + 2\eta_3 \sin^2 \alpha + \eta_2) \tilde{V}_j^T + (2\chi_3 \sin \alpha - \chi_2) \tilde{\mathcal{O}}_i - 6b_2 \tilde{\mathcal{O}}_j \\ & + 6b_3 \sin \alpha (\tilde{\mathcal{O}}_{i+1} + \tilde{\mathcal{O}}_{i-1}) - \chi_3 \sin^2 \alpha (\tilde{V}_{j+1}^T + \tilde{V}_{j-1}^T) - \chi_3 \sin \alpha \cos \alpha (\tilde{\omega}_{j+1}^T - \tilde{\omega}_{j-1}^T) \\ & + \chi_2 \tilde{V}_j^0 = 0 \end{aligned}$$

$$\begin{aligned} \theta) & -(2\xi_3 \sin^2 \alpha + 2\eta_3 \cos^2 \alpha + \xi_2) \tilde{\omega}_j^T + \xi_2 \tilde{\omega}_j^0 + \xi_3 \sin \alpha \cos \alpha (\tilde{V}_{j+1}^T - \tilde{V}_{j-1}^T) \\ & - \xi_3 \sin^2 \alpha (\tilde{\omega}_{j+1}^T + \tilde{\omega}_{j-1}^T) + \left( \frac{1}{m\pi} \right)^2 (2\xi_3 \sin \alpha - \epsilon_2) \tilde{U}_i - \left( \frac{1}{m\pi} \right)^2 \mu_2 \cdot \tilde{U}_j \\ & \left( \frac{1}{m\pi} \right)^2 \mu_3 \sin \alpha (\tilde{U}_{i+1} + \tilde{U}_{i-1}) - 6b_3 \cos \alpha (\tilde{\mathcal{O}}_{i+1} - \tilde{\mathcal{O}}_{i-1}) + \chi_3 \sin \alpha \cos \alpha (\tilde{V}_{j+1}^T - \tilde{V}_{j-1}^T) \\ & - \chi_3 \cos^2 \alpha (\tilde{\omega}_{j+1}^T + \tilde{\omega}_{j-1}^T) = 0 \end{aligned}$$

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$$\begin{aligned} a_i &= \frac{K_i}{K_i^2}, \quad b_i = \frac{K_i}{K_i^2}, \quad \alpha_i = K_i \left( \frac{1}{K_i^2} + \nu \left( \frac{m\pi}{l} \right)^2 \right), \quad \beta_i = \frac{G_i}{K_i} + \frac{N_i K_i}{3} \left( \frac{m\pi}{l} \right)^2, \quad \gamma_i = \frac{G_i}{K_i} - \frac{N_i K_i}{3} \left( \frac{m\pi}{l} \right)^2 \\ \epsilon_i &= \left( \frac{m\pi}{l} \right)^2 \left( \frac{\nu N_i}{2} - \frac{G_i}{2} \right), \quad \mu_i = \left( \frac{m\pi}{l} \right)^2 \left( \frac{\nu N_i}{2} + \frac{G_i}{2} \right), \quad \zeta_i = \frac{N_i}{K_i} + \frac{G_i K_i}{3} \left( \frac{m\pi}{l} \right)^2, \quad \xi_i = \frac{N_i}{K_i} - \frac{G_i K_i}{3} \left( \frac{m\pi}{l} \right)^2, \\ \eta_i &= K_i \left( \frac{12}{K_i^3} + \frac{2}{K_i^2} \left( \frac{m\pi}{l} \right)^2 + \frac{K_i}{3} \left( \frac{m\pi}{l} \right)^4 \right), \quad \chi_i = K_i \left( \frac{12}{K_i^3} + \frac{2}{K_i^2} \left( \frac{m\pi}{l} \right)^2 - \frac{K_i}{3} \left( \frac{m\pi}{l} \right)^4 \right) \\ \mathcal{X}_i &= K_i \left( \frac{6}{K_i^2} + \left( \frac{m\pi}{l} \right)^2 \right) \quad (i=1, 2, 3) \end{aligned}$$

ii) 有限7-リ変換

各変形に、加えて次の有限7-リ変換を行う。

$$\begin{aligned} \tilde{\mathcal{O}}_j^0 &= \sum_{k=1}^n \tilde{\mathcal{O}}_k^0 \cos \frac{k\pi}{n} j & \tilde{\mathcal{O}}_j^T &= \sum_{k=1}^n \tilde{\mathcal{O}}_k^T \cos \frac{k\pi}{n} j \\ \tilde{V}_j^0 &= \sum_{k=1}^n \tilde{V}_k^0 \cos \frac{k\pi}{n} j & \tilde{V}_j^T &= \sum_{k=1}^n \tilde{V}_k^T \cos \frac{k\pi}{n} j \\ \tilde{\omega}_j^0 &= \sum_{k=1}^n \tilde{\omega}_k^0 \sin \frac{k\pi}{n} j & \tilde{\omega}_j^T &= \sum_{k=1}^n \tilde{\omega}_k^T \sin \frac{k\pi}{n} j \\ \tilde{U}_j^0 &= \sum_{k=1}^n \tilde{U}_k^0 \sin \frac{k\pi}{n} j & \tilde{U}_j^T &= \sum_{k=1}^n \tilde{U}_k^T \sin \frac{k\pi}{n} j \end{aligned}$$

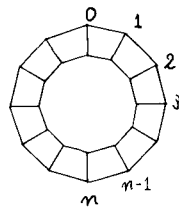


図. 4

その結果をマトリックスで表示すると

$$A \cdot U = P$$

$$\begin{aligned} U^T &= [\tilde{\mathcal{O}}_k^0 \quad \tilde{V}_k^0 \quad \tilde{\omega}_k^0 \quad \tilde{U}_k^0 \quad \tilde{\mathcal{O}}_k^T \quad \tilde{V}_k^T \quad \tilde{\omega}_k^T \quad \tilde{U}_k^T] \\ P^T &= [0 \quad 0 \quad \tilde{P} \quad 0 \quad 0 \quad 0 \quad 0 \quad 0] \end{aligned}$$

A は (8×8) のマトリックスで、各要素は  $K_{ij}$  で表示されたと同じである。

$$\begin{aligned} K_{11} &= 4(2a_1 + a_2 + a_3 \cos \frac{k\pi}{n}), \quad K_{12} = -(2a_1 \sin \alpha + a_2 + 12b_1 \sin \alpha \cos \frac{k\pi}{n}), \quad K_{13} = \\ & -12b_1 \cos \alpha \sin \frac{k\pi}{n}, \quad K_{14} = K_{17} = K_{18} = 0, \quad K_{15} = 2a_2, \quad K_{16} = 6b_2, \quad K_{21} = 0, \\ K_{22} &= -2\mu_1 \cos \alpha \sin \frac{k\pi}{n}, \quad K_{23} = -(2\epsilon_1 \sin \alpha + \epsilon_2 + 2\mu_1 \sin \alpha \cos \frac{k\pi}{n}), \quad K_{25} = K_{26} = 0, \\ K_{24} &= (2\beta_1 + \beta_2 - 2\gamma_1 \cos \frac{k\pi}{n}), \quad K_{27} = \mu_2, \quad K_{28} = -\gamma_2, \quad K_{31} = -(2\chi_1 \sin \alpha + \chi_2 + 12b_3 \sin \alpha \cos \frac{k\pi}{n}) \\ & \times (\cos \frac{k\pi}{n}), \quad K_{32} = (2\xi_1 \cos^2 \alpha + 2\eta_1 \sin^2 \alpha + \eta_2) + 2(\chi_1 \sin^2 \alpha - \xi_1 \cos^2 \alpha) \cos \frac{k\pi}{n}, \quad K_{33} = 2(\chi_1 + \\ & \xi_1) \sin \alpha \cos \alpha \sin \frac{k\pi}{n}, \quad K_{34} = -\left( \frac{\epsilon}{m\pi} \right)^2 \mu_1 \cos \alpha \sin \frac{k\pi}{n}, \quad K_{35} = -6b_2, \quad K_{36} = -\chi_2, \quad K_{37} = K_{38} = 0 \\ K_{41} &= -12b_1 \cos \alpha \sin \frac{k\pi}{n}, \quad K_{42} = 2(\xi_1 + \chi_1) \sin \alpha \cos \alpha \sin \frac{k\pi}{n}, \quad K_{43} = (2\xi_1 \sin^2 \alpha + 2\eta_1 \cos^2 \alpha + \xi_2) + \end{aligned}$$

$$\begin{aligned}
& + 2(\xi_1 \sin^2 \alpha - \eta_1 \cos^2 \alpha) \cos \frac{k\pi}{n}, K_{4,4} = -\left(\frac{\rho}{m\pi}\right)^2 (2E_1 \sin \alpha + E_2 + 2\mu_1 \sin \alpha \cos \frac{k\pi}{n}), K_{4,5} = K_{4,6} = \\
& K_{4,7} = -\xi_2, K_{4,8} = -\left(\frac{\rho}{m\pi}\right)^2 \mu_2, K_{5,1} = -2\alpha_2, K_{5,2} = 6b_2, K_{5,3} = K_{5,4} = 0, K_{5,5} = -4(2\alpha_3 + \alpha_2 + \alpha_3) \\
& \times \cos \frac{k\pi}{n}, K_{5,6} = (2\alpha_3 \sin \alpha - \alpha_2 + 12b_3 \sin \alpha \cos \frac{k\pi}{n}), K_{5,7} = 12b_3 \cos \alpha \sin \frac{k\pi}{n}, K_{5,8} = 0, \\
& K_{6,1} = K_{6,2} = 0, K_{6,3} = \mu_2, K_{6,4} = \gamma_2, K_{6,5} = 0, K_{6,6} = 2\mu_3 \cos \alpha \sin \frac{k\pi}{n}, K_{6,7} = 2E_3 \sin \alpha - E_2 \\
& + 2\mu_3 \sin \alpha \cos \frac{k\pi}{n}, K_{6,8} = -(2b_3 + 3z - 2\gamma_3 \cos \frac{k\pi}{n}), K_{7,1} = -6b_2, K_{7,2} = \gamma_2, K_{7,3} = K_{7,4} = 0 \\
& K_{7,5} = 2\alpha_3 \sin \alpha - \alpha_2 + 12b_3 \sin \alpha \cos \frac{k\pi}{n}, K_{7,6} = -(2\xi_3 \cos^2 \alpha + 2\eta_3 \sin^2 \alpha + \gamma_2) - 2(\alpha_3 \sin^2 \alpha - \\
& \xi_3 \cos^2 \alpha) \cos \frac{k\pi}{n}, K_{7,7} = -2(\alpha_3 + \xi_3) \sin \alpha \cos \alpha \sin \frac{k\pi}{n}, K_{7,8} = 2\left(\frac{\rho}{m\pi}\right)^2 \mu_3 \cos \alpha \sin \frac{k\pi}{n}, \\
& K_{8,1} = K_{8,2} = 0, K_{8,3} = \xi_2, K_{8,4} = -\left(\frac{\rho}{m\pi}\right)^2 \mu_2, K_{8,5} = 12b_3 \cos \alpha \sin \frac{k\pi}{n}, K_{8,6} = -2(\xi_3 + \alpha_3) \\
& \times \sin \alpha \cos \alpha \sin \frac{k\pi}{n}, K_{8,7} = -(2\xi_3 \sin^2 \alpha + 2\eta_3 \cos^2 \alpha + \xi_2) - 2(\xi_3 \sin^2 \alpha - \alpha_3 \cos^2 \alpha) \cos \frac{k\pi}{n} \\
& K_{8,8} = \left(\frac{\rho}{m\pi}\right)^2 (2E_3 \sin \alpha - E_2 + 2\mu_3 \sin \alpha \cos \frac{k\pi}{n})
\end{aligned}$$

又、荷重項は

$$\tilde{P}_j = \sum_{k=1}^n \tilde{R}_k \cdot \cos \frac{k\pi}{n} j, \quad C_k[\tilde{P}_j] = \sum_{j=1}^{n-1} \tilde{P}_j \cos \frac{k\pi}{n} j \quad (2)$$

但し、

$$\begin{aligned}
\tilde{R}_0 &= \frac{1}{n} \{ C_0[\tilde{P}_j] + \frac{1}{2} \tilde{P}_{(n)} + \frac{1}{2} \tilde{P}_{(0)} \} \\
\tilde{R}_k &= \frac{1}{n} \{ 2C_k[\tilde{P}_j] + \tilde{P}_{(n)}(-1)^k + \tilde{P}_{(0)} \} \\
\tilde{R}_n &= \frac{1}{n} \{ C_n[\tilde{P}_j] + \frac{1}{2} \tilde{P}_{(n)}(-1)^n + \frac{1}{2} \tilde{P}_{(0)} \}
\end{aligned}$$

今、考えている荷重状態では、 $\tilde{P}_{(n)} = \tilde{P}_{(0)} = 1$ ,  $C_0[\tilde{P}_j] = C_k[\tilde{P}_j] = C_n[\tilde{P}_j] = 0$  とする。  
すなわち、

$$\tilde{P}_j = \frac{1}{n} + \frac{1}{n} \sum_{k=1}^{n-1} [1 + (-1)^k] \cos \frac{k\pi}{n} j + \frac{1}{2n} [1 + (-1)^n] \cos \frac{n\pi}{n} j$$

第一項は、全節点に一定線荷重が作用する場合に相当する。従って、1, 2, 3項の荷重状態との重ね合わせに  
よって、図1に示す荷重状態に相当する変形を求めると出来る。

iii) 逆変換

$A \cdot U = P$  によって得られた変形を、逆変換する。

$$\theta_j^0 = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{\theta}_k^0 \sin \frac{m\pi x}{\ell} \cos \frac{k\pi}{n} j$$

$$U_j^0 = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{U}_k^0 \sin \frac{m\pi x}{\ell} \cos \frac{k\pi}{n} j$$

$$W_j^0 = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{W}_k^0 \sin \frac{m\pi x}{\ell} \sin \frac{k\pi}{n} j$$

$$U_j^I = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{U}_k^I \sin \frac{m\pi x}{\ell} \sin \frac{k\pi}{n} j$$

$$\theta_j^I = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{\theta}_k^I \sin \frac{m\pi x}{\ell} \cos \frac{k\pi}{n} j$$

$$U_j^I = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{U}_k^I \sin \frac{m\pi x}{\ell} \cos \frac{k\pi}{n} j$$

$$W_j^I = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{W}_k^I \sin \frac{m\pi x}{\ell} \sin \frac{k\pi}{n} j$$

$$U_j^I = \frac{2}{\ell} \sum_{m=1}^{\infty} \sum_{k=1}^n \tilde{U}_k^I \sin \frac{m\pi x}{\ell} \sin \frac{k\pi}{n} j$$

尚、数値計算は当日会場発表する予定である。

#### 4. 参考文献

i) 能町, 堀, 折板によるマルチセル大口径鋼管杭の応力解析について, 土木学会第28回年次学術講演会  
講演集第I部, S. 48, 10.

ii) Sumio. G. Nomachi; On a Stress Analysis of Grid Plate by Finite Fourier Transforms  
Concerning Finite Integration, Proceedings of the Sixteenth Japan National  
Congress for Applied Mechanics, 1966.