INVESTIGATION ON WAKE GALLOPING OF CABLE WITH SPIRAL PROTUBERANCES BY PRESSURE MEASUREMENTS

| 京都大学 | 学生員 | \bigcirc Tung Quang Do | 京都大学 | 正会員 | 八木 知己 |
|------|-----|--------------------------|------|-----|-------|
| 京都大学 | 正会員 | 野口 恭平 | 京都大学 | | 福島 温樹 |
| 京都大学 | | Thu Minh Dao | 京都大学 | 学生員 | 桑原 彰吾 |

1. INTRODUCTION

Wake-induced vibration is a major concern for cables arranged in parallel. There are several possible wake-induced vibration responses depending on the arrangement of the cables. This study focuses on wake galloping (or vertical flutter) of 2 relatively close cables. A previous study has shown that a cable with 12 spiral protuberances at winding angle of 27° (hereinafter called spiral cable) could partially suppress wake induced vibration in parallel cables¹). However, the suppression mechanism is yet to be investigated. This study discusses wake galloping characteristics of spiral cable at high reduced wind velocity based on surface pressure obtained from wind tunnel tests.



Fig. 1 Parallel cables arrangement in crossflow Pressure ring

Fig. 2 Spiral cable model

2. WAKE GALLOPING RESPONSES OF PARRALEL CABLES AT HIGH WIND VELOCITY

Fig. 1 shows the arrangement and coordinate systems for 2 parallel cables. (*X*,*Y*) and (*S*,*W*) denote the distance of the downstream cable from the upstream cable in wind-off (initial condition) and wind-on condition respectively; θ represents the angular position of the pressure taps; η and ξ are the displacements in vertical and horizontal direction. The main focus of this study is placed on (*X*/*D*,*Y*/*D*)=(3.0, 0.5) which is equivalent to (*S*/*D*,*W*/*D*)=(3.0, 0.3) and at reduced wind velocity of *U*/*fD*>100 (corresponding to *Re*>2×10⁴ – the supercritical region for a single spiral cable¹) in view of practical engineering applications as well as minimizing the effect of Reynolds dependency. In this particular arrangement where both cables were smooth circular cylinder, denoted by Circular (upstream) – Circular (downstream) in Fig. 3, large amplitude vertical flutter was observed in the downstream cable at *U*/*fD*>100 (the shaded areas in Fig. 3). A downstream spiral cable was able to suppress the vibration at high reduced wind velocity. It seemed that the configuration of the downstream cable strongly dictated its stability regardless of the configuration of the upstream one as there was no significant $\frac{\Theta}{\Theta}$

Flutter analysis revealed that the vibration observed in Circular – Circular was driven by unsteady aerodynamic derivative H_1^* which is defined in Eq. (1)²⁾.

$$H_1^* = -\frac{L_{\eta 0} \sin \psi_{L\eta}}{\rho (0.5D)^2 \omega^2 \eta_0}$$
(1)

where: ρ is the air density [kg/m³]; *D* is the diameter of the cable [m]; ω is the angular frequency of the forced vibration [rad/s]; $L_{\eta 0}$ is the amplitude of unsteady lift force per unit length [N/m]; $\psi_{L\eta}$ is the phase lag between vertical displacement and the unsteady





lift force; η_0 is the amplitude of the vertical motion in forced vibration. $H_1^* > 0$ describes negative aerodynamic damping in the vertical direction and thus, indicates galloping. Additionally, $H_1^* > 0$ was shown to have a major role in the 2DOF flutter with dominant vertical vibration of Circular – Circular seen in Fig. 3²).

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3. WAKE GALLOPING CHARACTERISTICS OF PARALLEL CABLES

In this section, wake galloping suppression mechanism of spiral cable is discussed based on static and unsteady pressure in (S/D, W/D)=(3.0, 0.3) at U/fD=183 (*Re*=4.2×10⁴).

Fig. 4 shows the distribution of surface pressure coefficient C_P on the circumference of the downstream cable. The high negative pressure at $0^{\circ} \le \theta \le 120^{\circ}$ in Circular – Circular originated from gap flow³), which is a high-speed flow in the gap between two parallel cables in certain arrangements. Gap flow generated non-zero lift force on the downstream cable and directed it toward the upstream one. In the case of downstream spiral cable, this negative pressure region was significantly reduced. Thus, the spiral protuberances on the downstream cable must have somehow suppressed the gap flow.

Fig. 5 shows H_1^* calculated at each pressure tap on the circumference of the downstream cable. In the case of downstream smooth cable, $H_1^*>0$ was found at the outer region of the cable ($180^\circ \le \theta \le 360^\circ$). H_1^* was maximum at $255^\circ \le \theta \le 270^\circ$. The majority of $H_1^*>0$ did not come from the gap flow region but the outer region. As shown in Eq. (1), H_1^* was decided by $L_{\eta 0}$ (integrated from C_P) and $\psi_{L\eta}$. At the outer region, although negative pressure was not as significant as in the gap flow region, there existed a remarkable phase lag. Time series $C_P(t)$ of the downstream smooth cable is presented in Fig. 6a. High negative pressure (denoted by the dark blue region) at the inner region did not persistently exist throughout the whole period T but rather switched intermittently between the inner region ($0^\circ \le \theta \le 90^\circ$) and the outer region ($255^\circ \le \theta \le 270^\circ$). It seemed that the gap flow periodically reduced its strength as it switched to the outer flow. This flow switching phenomenon probably caused the significant phase lag at the outer region.

Such switching phenomenon, however, was not observed in the case of downstream spiral cable (Fig. 6b and 6c). The spiral protuberances must have suppressed the gap flow so that no switching occurred. Consequently, not only did negative pressure get minimized but the phase lag was also reduced and H_1^* decreased. Fig. 5 reaffirmed the action of the downstream spiral cable was consistent regardless of the configuration of the upstream cable as H_1^* profiles in Circular - Spiral and Spiral - Spiral are very close in shape. It should be noted that there existed variation in the unsteady pressure field along the span of the spiral cable. However, wake galloping and its suppression mechanism could be explained in the view of 2dimensional pressure characteristics because the sign of H_1^* remained consistent spanwise.

4. CONCLUSIONS

1) Gap flow did not directly drive wake galloping but the flow switching phenomenon between gap flow and outer flow.



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Fig. 4 C_P in circumferential direction at U/fD=183





of the forced vibration $G_{i}^{(i)} = 1 + G_{i}^{(i)} +$

