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1. Introduction

The generation and dissipation of pore water pressure is very important in soil mechanics, geotechnical engineering, and geoenvironmental engineering. A lot of works have been done on the problem, but there still is a big gap between the requirement and the research, this may be due to the fact that the different materials have different responses[2]. In this paper we will discuss the effect of different constitutive models on the dissipation of pore water pressure by means of Biot's consolidation theory[1,3]. From the discussion we can see that even almost the same response in derivatoric stress vs strain in drainage but a great difference exists in pore water pressure.

2. Components of Biot's theory incorporating with general constitutive models Following six concepts form generalized Biot's theory[1,3]:

• Equilibrium equation

$$\frac{\partial \sigma_{ij}}{\partial x_i} + b_i = 0$$

• Constitutive laws of solid skeleton

$$d\sigma'_{ij} = D_{ijkl}d\varepsilon_{kl}$$

• Darcy's seepage law

$$q_i = K_{ij} \frac{\partial \psi}{\partial x_j}$$

• Geometrical equation

$$\Delta \varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial \Delta u_i}{\partial x_i} + \frac{\partial \Delta u_j}{\partial x_i} \right)$$

• Terzaghi's effective stress principle

$$\sigma_{ij} = \sigma'_{ij} + u\delta_{ij}$$

• Incompressibility of solid-water mixture

$$\frac{\partial \varepsilon_v}{\partial t} = \frac{\partial q_i}{\partial x_i}$$

3. Weak form and FEM discretization

The body forces on the skeleton are two components: 1) The effective weight b_i per unit volume. 2) Seepage force induced by $(-u_{i})$. So that the weak form for error-self-corrector mode is:

$$\int_{V} \{\delta(\Delta \varepsilon)\}^{T} [D]_{ep} \{\Delta \varepsilon\} dv - \int_{V} \{\delta(\frac{\partial \Delta \bar{u}_{i}}{\partial x_{i}})\}^{T} \{u\}^{t+\Delta t} dv - \int_{V} \{\delta(\Delta \bar{u})\}^{T} \{\Delta b\} dv
+ \int_{S_{\sigma}} \{\delta(\Delta \bar{u})\}^{T} \{n\} u^{t+\Delta t} ds - \int_{S_{\sigma}} \{\delta(\Delta \bar{u})\}^{T} \{\Delta \bar{T}\} ds
= -\int_{V} \{\delta(\Delta \varepsilon)\}^{T} \{\sigma\} dv + \int_{S_{\sigma}} \{\delta(\Delta \bar{u})\}^{T} \{\bar{T}\} ds - \int_{V} \{\delta(\Delta \bar{u})\}^{T} \{b\} dv$$

$$- \int_{V} \{\delta u\}^{T} \{\frac{\partial \Delta \bar{u}_{i}}{\partial x_{i}}\} dv = \frac{1}{r_{w}} \int_{t}^{t+\Delta t} [\int_{S_{\sigma}} \{\delta u\}^{T} \{K u\} ds] dt - \frac{1}{r_{w}} \int_{t}^{t+\Delta t} [\int_{V} \{\frac{\partial \delta u}{\partial x_{i}}\}^{T} \{K_{i} \frac{\partial u}{\partial x_{i}}\} dv] dt$$
(2)

For plane strain problems $(i=1\Rightarrow x;\ i=2\Rightarrow y;$ element domain V_m), the displacement \bar{u} and excess pore pressure u are taken to be the same shape function, 4-nodal isoparametric element. Crank-Nicolson mode $(\theta=\frac{1}{2})$ is adopted for time domain.

4. Constitutive Models of A Compacted Clay

4.1 Elastoplastic model

$$\{d\sigma\} = [D]_{ep} \{d\varepsilon\}, \qquad [D]_{ep} = [D] - \frac{[D] \{\frac{\partial g}{\partial \sigma}\} \{\frac{\partial f}{\partial \sigma}\}^T [D]}{\bar{A} + \{\frac{\partial f}{\partial \sigma}\}^T [D] \{\frac{\partial g}{\partial \sigma}\}} = [D] - \frac{G[X]}{\frac{\bar{A}}{G} + \Phi}$$
(3)

Here [D] is elastic matrix; G: shear elastic modulus; \bar{A} : plastic hardening modulus; g: plastic potential function; f: yielding function; $\Phi = \{\frac{\partial f}{\partial \sigma}\}^T [D] \{\frac{\partial g}{\partial \sigma}\}/G$. Associated flow rule is used. The yield function f and hardening function H of a compacted clay are

$$f = \left(\frac{p' - H}{1.206H}\right)^2 + \left(\frac{q}{1.46H}\right)^2 - 1 = 0 \qquad H = \frac{P_a}{2.206} \left(\frac{\epsilon_v^p - 0.06\bar{\epsilon}^p}{0.0166}\right)^{\frac{1}{0.5}} + \frac{P_r}{2.206}$$
(4)

 $\begin{cases} H \geq H_{max} & : \text{loading} \\ H < H_{max} & : \text{unloading/reloading} \end{cases}$ (5) $P_{\tau} \text{ is the overconsolidation pressure, } P_{a} \text{ is the atmospheric pressure. And elastic bulk modulus}$ B = 157.8p', and $G = 103.33P_a \left(\frac{\sigma_3}{P_s}\right)^0$

4.2 Duncan-Change nonlinear elasticity[3]

$$E_t = 20.2P_a \left(8 + \frac{\sigma_3}{P_a}\right)^{1.0944} \left[1 - \frac{0.7156(1 - \sin 30.7^\circ)(\sigma_1 - \sigma_3)}{0.4P_a \cos 30.7^\circ + 2\sigma_3 \sin 30.7^\circ}\right]^2$$
(6)

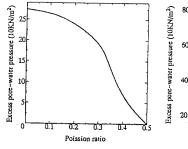
$$E_{ur} = 347.1 P_a \left(\frac{\sigma_3}{P_a}\right)^{0.8304} \qquad B = 14.64 P_a \left(8 + \frac{\sigma_3}{P_a}\right)^{1.08} \tag{7}$$

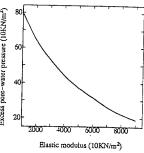
The permeability of this compacted clay is $K = 1.6 \times 10^{-7} m/s$.

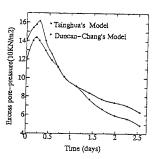
5. Cases study

One-dimensional consolidation

• Two-dimensional consolidation







6. Conclusions

- · Dilatancy has a vital effect on excess pore water pressure. It is dangerous to ignore the positive dilatancy when the generation and dissipation of pore water pressure is considered.
 - Loading level will affect the Mandel-Cryer effect.

References

[1] 市川康明 (1990), 地盤力学における有限要素法入門, 日科技連出版社. [2] Jian-Guo Wang, T. Ito, and Yasuaki Ichikawa(1995) Effect of dilatancy on pore water pressure in porous medium, 第26回岩盤力学 に関するシンポジウム講演論文集(東京, 1995 年 1 月 19 日 — 20 日), pp.480-484 [3] Wang, J.-G.(1994), Instructions on COND2.FORT(2-DIM), Ichikawa's Lab., Dept. of Geotechnical and Environmental Engineering, Nagoya University